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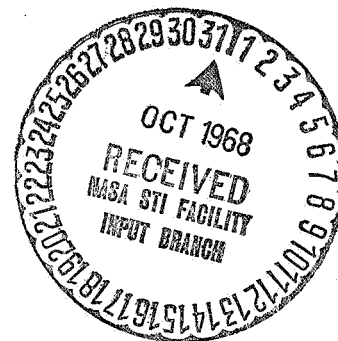
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VELOCITY FIELD EXCITED BY WING VIBRATIONS  
PROPAGATING ALONG ITS SURFACE  
WITH SUPERSONIC VELOCITY

by

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VELOCITY FIELD EXCITED BY WING VIBRATIONS  
PROPAGATING ALONG ITS SURFACE  
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by E. A. Krasil'shchikova

SUMMARY

The problem discussed in the present paper results in defining a function which determines the amplitude and the initial phase of oscillations at each oscillating point of the surface of the wing.

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1. Let us consider a rectilinear forward motion of a wing with constant velocity  $\underline{u}$  inside a boundless volume of ideal compressible medium. Beginning from a certain moment of time  $t_0$ , small oscillations propagate along the wing's elastic surface with supersonic velocity  $\underline{v}$ . The normal velocity component, due to the basic motion, is given on both sides of the wing in the form

$$v_{0n} = -u\alpha, \quad (1)$$

where  $\alpha$  is the angle of attack of streamlined surface's elements. The normal velocity component due to vibrations is given in the form

$$v_{\Delta n} = A_{\Delta}, \quad (2)$$

Where  $A_{\Delta}$  is a function of time and points of streamlined surface. Functions and  $A_{\Delta}$  are small and may be arbitrary integrable functions of their arguments.

Considering that the medium is weakly disturbed, we shall consider the problem of determination of the field of velocities in linearized setup [1, 2]. We shall assume the medium's motion as irrotational and taking place in the absence of external forces. Let us take a mobile axis system of coordinates  $Oxyz$ , invariably linked with the moving wing (Fig.1). The axis  $Oz$  is directed perpendicularly to the plane of the drawing.

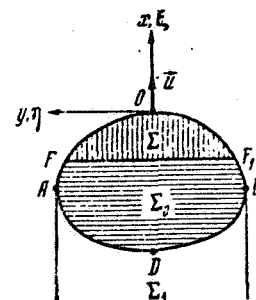


Fig.1

The potential of velocity  $\phi$  satisfies the equation

$$(u^2 - a^2)\phi_{xx} - a^2\phi_{yy} - a^2\phi_{zz} - 2u\phi_{xt} + \phi_{tt} = 0, \quad (3)$$

where  $a$  is the speed of sound in an unperturbed medium, and the boundary conditions in the plane  $xOy$ .

In the region  $\Sigma_1$ , i.e. the projection of the wing on the plane  $xOy$  ahead of the front of vibrations' propagation (line  $FF_1$  in Fig.1), the derivative is

$$\phi_z = -ua(x, y) = A_0(x, y). \quad (4)$$

In the region  $\Sigma$  i.e., the projection of the wing on the plane  $xOy$  behind the front  $FF_1$ , the derivative is

$$\phi_z = A_0(x, y) + A_\Delta(x, y, t) = A(x, y, t). \quad (5)$$

In the region  $\Sigma_1$ , i. e., the projection of the vortical shroud on the plane  $xOy$ , we have

$$\phi_t - u\phi_x = 0. \quad (6)$$

The potential  $\phi$  is zero everywhere in the plane  $xOy$  outside the region  $\Sigma_0 + \Sigma + \Sigma_1$ ; thus,

$$\phi = 0. \quad (7)$$

Besides, the Chaplygin-Zhukovskiy's principle must be observed at the trailing edge of the wing at any moment of time.

Therefore, the problem consists in the following: to find in the half-space  $z \geq 0$  a function  $\phi(x, y, z, t)$  that would satisfy Eq.(3), the boundary conditions (4)-(5), given in regions with mobile boundary, and the boundary conditions (6) - (7), given in regions with fixed boundaries. The solution of the problem in the half-space  $z < 0$  will be found from the condition  $\phi(x, y, -z, t) = -\phi(x, y, z, t)$ .

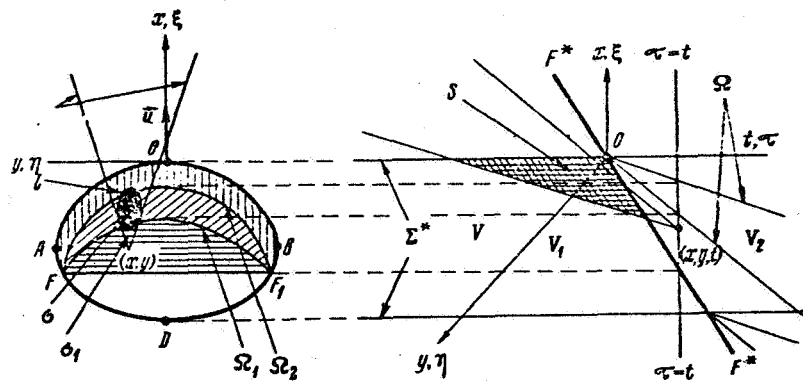


Fig. 2

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2. For the solution of this problem we shall apply the method developed in papers [2, 4]. Let us turn to space  $xyt$  and consider in it a region  $V$ , inside of which the derivatives  $\phi_z$  with respect to flowing around condition are given. Region  $V$  is bounded by surface  $\Sigma^*$ . This surface  $\Sigma^*$  is a cylindrical surface with generatrices parallel to time axis  $Ot$ , and a directrix representing the contour AOB $\bar{D}$ , which is given by the equation (Fig.2)

$$\eta = \psi(\xi)$$

Let the line  $FF_1$ , - the propagation front of vibrations along the wing's surface be displaced according to the law:  $x = f(t)$ , whereupon  $f'(t) = v$ . The surface  $F^*$ , given by equation  $\xi = f(\tau)$ , divides the region  $V$  into two parts  $V_1$  and  $V_2$  with different values of the derivative  $\phi_z$ , according to conditions (4) and (5). To the left of the surface  $F^*$ , in the region  $V_1$ , the derivative is  $\phi_z = A_0$ , and to the right, in the region  $V_2$ , it is  $\phi_z = A$  (Fig.2).

Let us take the solution of Eq.(3) in the form [3]

$$\varphi(x, y, z, t) = \frac{u^2 - a^2}{2\pi} \iint_{S(x, y, z, t)} \frac{\varphi_z(\xi, \eta, 0, t - \frac{u(x - \xi) + ar}{u^2 - a^2}) dS}{V(u^2 - a^2)r^2 + [a(x - \xi) - ur]^2 + a^2k^2(y - \eta)^2}, \quad (8)$$

$$r = \sqrt{(x - \xi)^2 - k^2(y - \eta)^2 - k^2z^2}, \quad k = \sqrt{u^2/a^2 - 1},$$

where the integration region  $S$  is the surface of a hyperboloid determined by the equation

$$(x + \xi)^2 + (y - \eta)^2 + z^2 + 2u(x - \xi)(t - \tau) + (u^2 - a^2)(t - \tau)^2 = 0 \quad (9)$$

and the inequality  $\tau < t$ . We shall pass in formula (8) from the surface integral to double integrals with plane region of integration in the plane  $xOy$ . Let us consider a supersonic wing velocity  $u > a$ ; then we shall obtain

$$\begin{aligned} \varphi(x, y, z, t) = & -\frac{1}{2\pi} \iint_{S^*(x, y, z)} \frac{\varphi_z(\xi, \eta, 0, \tau_1)}{r} d\xi d\eta - \\ & -\frac{1}{2\pi} \iint_{S^*(x, y, z)} \frac{\varphi_z(\xi, \eta, 0, \tau_2)}{r} d\xi d\eta. \end{aligned} \quad (10)$$

$$\tau_1 = t + [u(x - \xi) + ar]/(u^2 - a^2), \quad \tau_2 = t + [u(x - \xi) - ar]/(u^2 - a^2).$$

Region  $S^*$  is bounded from above by a Mach wave, and from below by Mach hyperbola, or, as  $z = 0$ , by Mach lines.

When constructing the solution, an essential role is played by the line of intersection of surfaces  $S$  and  $F^*$ . We shall denote the projection of this line on the plane  $xOy$  by  $l$ . Curve  $l$  subdivides the plane region  $S^*$  into parts with different values of the derivative  $\phi_z$ . If the propagation front of vibrations along the wing's surface constitutes a straight line  $FF_1$ , shifting oppositely to the motion of the wing with constant velocity  $v > u + a$ , the surface  $F^*$  is a plane determined by the equation  $\xi + v\tau = 0$ , and the curve  $l$

is an ellipse determined by the equation

$$v^2(x - \xi)^2 + v^2(y - \eta)^2 + v^2z^2 + 2uv(x - \xi)(vt + \xi) + (u^2 - a^2)(vt + \xi)^2 = 0. \quad (11)$$

Note that ellipse  $l$  is always written into the Mach hyperbola, or at  $z = 0$ , into the angle formed by Mach lines.

Assuming that the condition

$$|u\psi'(\xi) / \sqrt{\psi'^2(\xi) + 1}| \geq a,$$

is fulfilled on the contour AOB $\bar{D}$ , we shall represent the solution (10) in the form

$$\varphi(x, y, z, t) = \varphi_0(x, y, z) - \frac{1}{2\pi} \iint_{\sigma(x, y, z, t) + \sigma_1(x, y, z, t)} \frac{A_\Delta(\xi, \eta, \tau_1)}{r} d\xi d\eta - \frac{1}{2\pi} \iint_{\sigma_1(x, y, z, t)} \frac{A_\Delta(\xi, \eta, \tau_2)}{r} d\xi d\eta, \quad (12)$$

where the integration region  $\sigma$  is the part of the region  $S^*$  that was found to be inside the ellipse  $l$  at the moment of time  $t$ , and the region  $\sigma_1$  is the part of the region  $S^*$  situated outside the ellipse  $l$ , below the arc of ellipse  $K_1LK_2$ . The points  $K_1$  and  $K_2$  are the points where ellipse  $l$  is tangent to the Mach hyperbola (Fig.3. In formula (12, function  $\phi_0$  is the solution of the well known problem of settled supersonic gas flow past the considered wing [5].

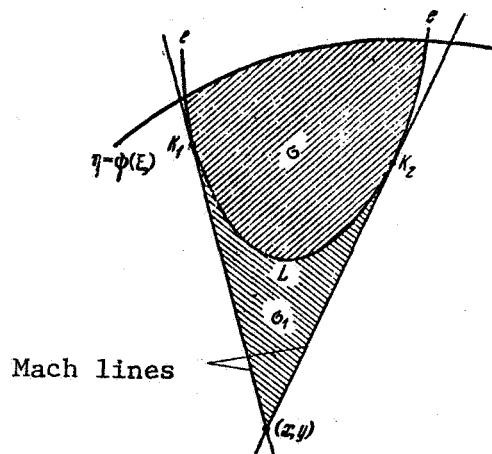


Fig.3

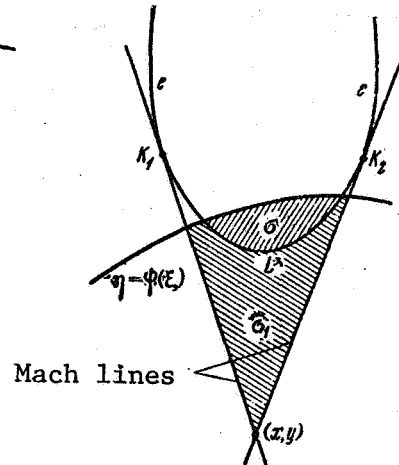


Fig.4

3. The analytical form of problem's solution depends on the mutual disposition of ellipse  $l$  and of wing contour AOB $\bar{D}$ , which determines the integration regions  $\sigma$  and  $\sigma_1$  in the solution (12).

Let us examine the moment of time  $t$ , belonging to the interval

$$0 = t_0 \leq t \leq t_1 = d/v,$$

where  $d$  is the length of segment OD. Let us draw a plane  $\tau = t$  in the space  $xyt$ . The projection of the line of plane's  $\tau = t$  intersection with the plane  $F^*$  on the plane  $xOy$  will be denoted by  $FF_1$  (Fig.2). Let us consider in the space  $xyt$  the family of cones determined by the equation

$$(X - \xi)^2 + (Y - \eta)^2 + 2u(X - \xi)(T - \tau) + (u^2 - a^2)(T - \tau)^2 = 0$$

and the inequality  $\tau > T$ , with summits on the line of intersection of Planes  $F^*$  and the surface  $\Sigma^*$ . We shall denote the envelope of this family by  $\Omega$ . The projections of the line of intersection of the plane  $\tau = t$  with the envelope  $\Omega$  on the plane  $xOy$  will be denoted by  $\Omega_1$  and  $\Omega_2$ . Let us find the equations of curves  $\Omega_1$  and  $\Omega_2$  in parametric form

$$\begin{aligned} v^2(x^* - \xi)^2 + v^2[\psi(x^*) - \eta]^2 - 2uv(x^* - \xi)(x^* + vt) + \\ + (u^2 - a^2)(x^* + vt)^2 = 0, \\ v^2(x^* - \xi) + v^2[\psi(x^*) - \eta]\psi'(x^*) - uv(x^* - \xi) - \\ - uv(x^* + vt) + (u^2 - a^2)(x^* + vt) = 0, \end{aligned} \quad (13)$$

where  $x^*$  is a parameter.

The lines  $FF_1$ ,  $\Omega_1$ ,  $\Omega_2$  divide the regions with different analytical character of problem's solution. To the region comprised between the straight line  $FF_1$  and the curve  $\Omega_1$ , responds the solution (12) in which the integration region  $\sigma$  is part of  $S^*$  bounded by ellipse  $l$ , located fully inside  $S^*$  (Fig.2). To the region comprised between curves  $\Omega_1$  and  $\Omega_2$ , responds the solution (12), in which  $\sigma$  is part of  $S^*$ , cut by ellipse  $l$ , crossing the contour of the wing. At the same time, the points of tangency  $K_1$  and  $K_2$  may be located inside (Fig.3), as well as outside of the wing (Fig.4). To the region comprised between curve  $\Omega_2$  and the Mach wave, responds the solution (12) when the integration in both integrals extends over the entire region  $S^*$ , that is, region  $\sigma$  is absent, while region  $\sigma_1$  coincides with the region  $S^*$ .

If the propagation velocity of vibrations along the streamlined surface satisfies the inequality  $v < u + a$ , curve  $l$  is a hyperbola.

4. In particular, when harmonic oscillations with frequency  $\omega$  propagate along the elastic surface of the wing, the function

$$A_\Delta(x, y, t) = A_1(x, y) \exp i\omega(t + a_1(x, y)) = \operatorname{Re} A_2(x, y) \exp i\omega t.$$

Utilizing the relation

$$\exp \frac{i\omega a}{u^2 - a^2} r + \exp - \frac{i\omega a}{u^2 - a^2} r = 2 \cos \frac{\omega a}{u^2 - a^2} r,$$

we shall represent the solution (12) in the form

$$\begin{aligned} \varphi = & \varphi_0 - \frac{1}{2\pi} \operatorname{Re} \exp(i\omega t) \exp\left(\frac{i\omega u}{u^2 - a^2} x\right) \times \\ & \times \iint_0 \frac{A_1(\xi, \eta)}{r} \exp(i\omega t) \exp\left(\frac{-i\omega u}{u^2 - a^2} \xi\right) \exp\left(\frac{i\omega a}{u^2 - a^2} r\right) d\xi d\eta - \\ & - \frac{1}{2\pi} \operatorname{Re} \exp(i\omega t) \exp\left(\frac{i\omega u}{u^2 - a^2} x\right) \times \\ & \times \iint_{a_1} \frac{A_2(\xi, \eta)}{r} \exp\left(-\frac{i\omega u}{u^2 - a^2} \xi\right) \cos \frac{\omega a}{u^2 - a^2} r d\xi d\eta, \end{aligned}$$

where the given function  $A_2$  determines the amplitude and the initial phase of oscillations at each oscillating point on the surface of the wing.

\*\*\*\* T H E E N D \*\*\*\*

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#### REFERENCES

- [1]. L. I. SEDOV. Metody podobiya i pazmernosti v mekhanike (Methods of Similitude and Dimensionality in Mechanics)., 1957.
- [2]. L. I. SEDOV. Ploskiye zadachi gidrodinamiki i aerodinamiki. (Plane Problems of Hydrodynamics and Aerodynamics), 1966.
- [3]. E. A. KRASIL'SHCHIKOVA. Dokl.AN SSSR, 117, 5, 1957
- [4]. E. A. KRASIL SHCHIKOVA. Arch. Mdch. Stosowanej, 2, 16, 1964.
- [5]. E. A. KRASIL SHCHIKOVA. Krylo konechnogo razmakha v szhimayemom potoke (Wing of finite span in compressible flow), 1952.

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